

NSF CyberTraining Winter Workshop

Introduction to Proper Orthogonal Decomposition

concepts, formulation and applications

Part I: Physics Simulations and POD Fundamentals

Ming-Cheng Cheng

Department of Electrical & Computer Engineering Clarkson University, Potsdam, NY 13699



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Applications of POD Simulation Methodology Physics Simulation of Domain Structures with Spatial Details

Physics Simulation of Domain Structures with Spatial Details

Numerical solution with detailed spatial solution in multidimensional structures:

Extremely time consuming. Why?

- Because of a large number of grids (nodes, or degrees of freedom) needed to solve the physical quantity numerically
- Usually, several 100 thousands or millions of nodes for a 3D problem
- Several 100 thousands or millions of coupled differential equations are needed



Mesh of a MOSFET

https://www.researchgate.net/figure/Mesh-of-a-MOSFET-with-a-polysilicon-gate_fig1_267688735



Applications of POD Simulation Methodology

Physics Simulation of Domain Structures with Spatial Details



- Heat Transfer Problems: Heat Transfer Equation
- Nanostructures and Materials: Schrödinger Equation Quantum Wave Equation (Quantum Eigenvalue problem), relevant to DFT simulation
- Electromagnetics and Photonics: Eigenvalue Problem or dynamic Wave Propagation
- Charge Carrier Transport in Semiconductor Devices: Electron/hole transport equations in (6D) phase space
- **Phonon Transport in nanostructures/nanodevices:** Phonon Boltzman Transport Equation (6D) phase space



Fluid dynamics: Navier Stokes Equations, conservation of energy, momentum, and mass

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Heat Transfer problems, for example, for semiconductor technology: semiconductor devices, integrated circuits, etc.

$$\rho C_{h} \frac{\partial T(\vec{r},t)}{\partial t} = \nabla \cdot k \nabla T(\vec{r},t) + P_{d}(\vec{r},t) \qquad \begin{array}{c} \text{Clarkson} \\ \text{UNIVERSITY} \\ \text{defy convention} \end{array}$$







Heat Transfer problems: CPUs, GPUs, etc.



Core 17	Core 13	Core 9	Core 4
Core 16	Core 12	Core 8	Core 3
Core 15	Core 11	Core 7	Core 2
Core 14	Core 10	Core 6	Core 1





Tesla Volta GV100 GPU with 13,440 cores, including FP32, FP64, INT32 and Tensor Cores.

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Quantum eigenvalue problem, Schrödinger equation

nanostructure, superlattice materials, density functional theory (DFT)



InGaAs/InAIAs Quantum Interband Cascade Laser (IEEE JQE Vol. 40, p. 1663, 2004)

Quantum Cascade Laser Power Supply



Qty 1+ \$5,599.00

Stock #15-957 \$5,599.00

Volume Pricing Request Quote



Si quantum-dot intermediate band solar cell

Nanotechnology, 24, 265401, 2013



Quantum Cascade Laser (QCL) Systems

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Quantum eigenvalue problem, Schrödinger equation

nanostructure, superlattice materials, density functional theory (DFT)

$$\nabla \cdot \left[-\frac{\hbar^2}{2m^*} \nabla \psi(\vec{r}) \right] + U(\vec{r})\psi(\vec{r}) = E\psi(\vec{r}),$$

Quantum Dot LEDs for TV and Cell phone Displays



Samsung: Quantum Dot LEDs

Color of Light Depends On Size of Quantum Dot

https://news.samsung.com/global/why-are-quantumdot-displays-so-good

Apple: Hybrid Quantum Dot LED and OLED Displays



https://www.patentlyapple.com/2019/10/apple-won-42-patentstoday-covering-hybrid-quantum-dot-displays-guis-supporting-3dar-models-a-health-study-more.html



Clarkson UNIVERSITY defy convention

Clarkson **Electromagnetics or photonics structures:** Dynamic wave equation, wave UNIVERSITY equation in frequency domain, EM eigenvalue problem)

Dynamic EM Wave equation: $\mu \epsilon \frac{\partial^2 \vec{E}(\vec{r},t)}{\partial t^2} + \mu \sigma \frac{\partial \vec{E}(\vec{r},t)}{\partial t} - \nabla^2 \vec{E}(\vec{r},t) = -\mu \frac{\partial \vec{j}(\vec{r},t)}{\partial t}$



https://en.wikipedia.org/wiki/Electromagnetic radiation







Diffraction of a Gaussian beam on a grating structure simulated using the FDTD (finite-difference timedomain) method.

(https://www.photond.com/products/omnis im/omnisim_applications 06.htm)

Electromagnetic Eigenvalue Problems

Electromagnetic Band Gap (EBG) Structures

$$\nabla^2 \vec{E}(\vec{r}) = -\left(\frac{\omega}{c}\right)^2 \epsilon_r \vec{E}(\vec{r})$$

Bloch Function: $\vec{E}(\vec{r}) = \vec{u}(\vec{r})e^{i\vec{k}\cdot\vec{r}}$



Photonic Crystals

Applications:

- Color selections
- Sensors
- Optical filters
- Photonic Crystal fibers
- Optical computing

https://www.tech-faq.com/what-are-photonic-crystals.html https://en.wikipedia.org/wiki/Photonic_crystal





Charge carrier transport in semiconductor devices (TCAD):

Boltzmann transport equation (BTE), hydrodynamic, energy transport model, drift-diffusion model



M. Shen, T. Zhou, MC Cheng, R. Fithen, *Computer Methods Appl. Mech & EAQ* Vol. 190, 2875-2891, February 16, 2001.

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How to develop a very efficient model (compact model) for a structure that requires the solution with spatial details?



Reduce the **numerical degrees of freedom (DoF)** by several orders of magnitude! **How? Projection or Mathematical Transformation**

Inner Product of vectors:

 $\vec{a} \cdot \vec{b} = |a||b|cos\theta = \sum_{i=1}^{M} a_i b_i$ = $a_x b_x + a_y b_y + a_z b_z$ (in real space)

L2 Norm:
$$|b| = \sqrt{\sum_{i=1}^{M} b_i^2}$$

Projection of \vec{a} onto \vec{b} :



Inner Product of functions:

$$f \cdot g = \int_{a}^{b} f(x)g(x)dx$$

L2 Norm:
$$|g| = \sqrt{\int_a^b g(x)^2 dx}$$

Projection of f onto g : $\frac{f \cdot g}{|g|}$

POD: Maximizing the average of (Projection)^2



• A popular projection, Fourier series (frequency domain projection)

$$f(t) = \sum_{n=1}^{M} a_n e^{jn\omega t} = \sum_{n=1}^{M} a_n \eta_n(t), \text{ or in space } f(x) = \sum_{n=1}^{M} a_n e^{jnkx} = \sum_{n=1}^{M} a_n \eta_n(x)$$

where the basis functions (or modes) $e^{jn\omega t}$ or e^{jnkx} are effective only for **periodic** functions; $\omega = \frac{2\pi}{T}$ and $k = 2\pi/\lambda$

 Spherical Harmonic Expansion: Legendre Polynomials



 Cylindrical harmonic expansion: Bessel Functions



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convention

https://opg.optica.org/oe/fulltext.cfm?uri=oe-18-24-25299&id=208252



 Instead of assuming the functions for the modes, a learning algorithm can be used to extract the modes.

The most commonly used projection-based learning methods:

- PCA: Principal component analysis
- SVD: Singular value decomposition
- The machine learning methods only provide statistical models to minimize the statistical variation in the prediction
 - Offer no guideline on physical principles in the governing equation for the spatial variation and/or dynamic evolution for the physical quantity.
 - Only work well for prediction of one-to-one correspondence relation
 - Not accurate for dynamic (initial value) problems
 - Poor performance for sudden change in physical quantity
 - Work well for interpolation but poorly for extrapolation



Effective projections of both the data and the governing equation onto the POD space

$$Q(\vec{r},t) = \sum_{j=1}^{M} a_j(t) \eta_j(\vec{r})$$

• POD finds a mode $\eta(\vec{r})$ that maximizes its mean square inner product (projection) with the solution data

$$\left| \left[\frac{Q \cdot \eta}{|\eta|} \right]^2 \right\rangle = \left| \left(\left(\int_{\Omega} Q(\vec{r}, t) \eta(\vec{r}) d\Omega \right)^2 \right) \right/ \int_{\Omega} \eta(\vec{r})^2 d\Omega \right|$$

 $\langle \rangle$ indicate the average of the data collected over numerical observations (or snapshots in time or over the samples) accounting for the parametric variations.

• This maximization process leads to an eigenvalue problem of 2-point correlation of data

$$Q(\vec{r},t) \otimes Q(\vec{r}',t) \eta(\vec{r}') d\Omega' = \lambda \eta(\vec{r}),$$



which offers a minimum least square error with a smallest number of modes (DoF) if the training is properly done (i.e., if the data quality is sufficient).

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Continuous eigenvalue problem

 $\int_{\Omega'} \langle Q(\vec{r},t) \otimes Q(\vec{r}',t) \rangle \eta(\vec{r}') d\Omega' = \lambda \eta(\vec{r}),$

Q data sets are collected from detailed numerical simulation (DNS)

 $\mathsf{Discretization}$ $\langle Q(\vec{r},t) \otimes Q(\vec{r}',t) \rangle = \frac{1}{N_s} \sum_{j=1}^{N_s} \vec{Q}(\vec{r},t_j) \otimes \vec{Q}(\vec{r}',t_j) = \frac{1}{N_s} \sum_{j=1}^{N_s} \vec{Q}(t_j) \vec{Q}^T(t_j)$ for each sample: $\vec{Q} \otimes \vec{Q} = \vec{Q} \vec{Q}^T = \begin{bmatrix} Q_1 \\ Q_2 \\ \vdots \\ Q_{N_r} \end{bmatrix} [Q_1 Q_2 \dots Q_{N_r}] = \begin{bmatrix} Q_1 Q_1 & \cdots & Q_1 Q_j & \cdots & Q_1 Q_{N_r} \\ \vdots & \ddots & \vdots & \ddots & \vdots \\ Q_i Q_1 & \cdots & Q_i Q_j & \cdots & Q_i Q_{N_r} \\ \vdots & \ddots & \vdots & \ddots & \vdots \\ Q_{N_r} Q_1 & \cdots & Q_{N_r} Q_j & \cdots & Q_{N_r} Q_{N_r} \end{bmatrix},$

A matrix dimension of $N_r \times N_r$ may be too large to manage in 3D problems for a dense matrix. The method of snapshots will be introduced later to minimize the computational effort



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Example, heat conduction in semiconductor Chips $Q(\vec{r},t) = T(\vec{r},t) = \sum_{j=1}^{M} a_j(t) \eta_j(\vec{r}); \quad a_j(t) \text{ needs to be determined}$ **Heat conduction equation:** $\frac{\partial \rho CT(\vec{r},t)}{\partial t} - \nabla \cdot k \nabla T(\vec{r},t) = P_d(\vec{r},t)$



$$\int_{\Omega} \left(\frac{\partial \rho CT(\vec{r},t)}{\partial t} - \nabla \cdot k \nabla T(\vec{r},t) = P_d(\vec{r},t) \right) \eta_i(\vec{r}) d\Omega$$

Using the following identities:

$$\nabla \cdot (\eta_i k \nabla T) = \nabla \eta_i \cdot k \nabla T + \eta_i \nabla \cdot k \nabla T$$

Gauss's Law:
$$\int_{\Omega} \nabla \cdot (\eta_i k \nabla T) d\Omega = \int_{S} \eta_i k \nabla T \cdot d\vec{S}$$

The projection leads to the weak form for the heat condition equation

$$\int_{\Omega} \left(\eta_i \frac{\partial \rho CT}{\partial t} + \nabla \eta_i \cdot k \nabla T \right) d\Omega = \int_{\Omega} \eta_i P_d \, d\Omega + \int_S \eta_i k \nabla T \cdot d\vec{S}$$

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Example, heat conduction in semiconductor Chips



• Galerkin projection (transformation) onto the the *i*th POD mode, $\eta_i(\vec{r})$

$$\int_{\Omega} \left(\eta_i \frac{\partial \rho CT}{\partial t} + \nabla \eta_i \cdot k \nabla T \right) d\Omega = \int_{\Omega} \eta_i P_d \, d\Omega + \int_S \eta_i k \nabla T \cdot d\vec{S}$$

• Using $T(\vec{r},t) = \sum_{j=1}^{M} a_j(t)\eta_j(\vec{r}) \Rightarrow M$ mode POD model, a set of *M*-dimensional ODEs $\sum_{j=1}^{M} c_{i,j} \frac{da_j}{dt} + \sum_{j=1}^{M} g_{i,j}a_j = P_{pod,i},$ $c_{i,j} = \int_{\Omega} \rho C \eta_i \eta_j \, d\Omega, \quad g_{i,j} = \int_{\Omega} k \nabla \eta_i \cdot \nabla \eta_j \, d\Omega, \quad P_{pod,i} = \int_{\Omega} \eta_i P_d(\vec{x}, t) \, d\Omega - \int_{\Gamma} \eta_i (-k \nabla T) \cdot d\vec{S}$



Procedure for Constructing the POD Simulation Model

- 1. Data Collection from direct numerical simulation (DNS) $\rightarrow Q(\vec{r}, t)$
- 2. Solving the 2-point correlation eigenvalue problem for $\lambda_j \& \eta_j(\vec{r})$. Observe the Eigenvalue spectrum to determine the number of modes, *M*, where λ_j represents the mean squared information captured by η_j Or estimate the least square error based on

$$Err_{LS,M} = \sqrt{\sum_{i=M+1}^{N_S} \lambda_i / \sum_{i=1}^{N_S} \lambda_i}$$

- 3. Project the governing equation onto the POD Space (accounting for physical principles) \rightarrow a set of *M* ODEs for a_j
- 4. Evaluate the model parameters (coefficients of the ODEs

Implementation of POD in physics simulation

• Solve the ODEs to obtain a_j



• Post processing: the solution $Q(\vec{r}, t) = \sum_{j=1}^{M} a_j(t) \eta_j(\vec{r})$



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